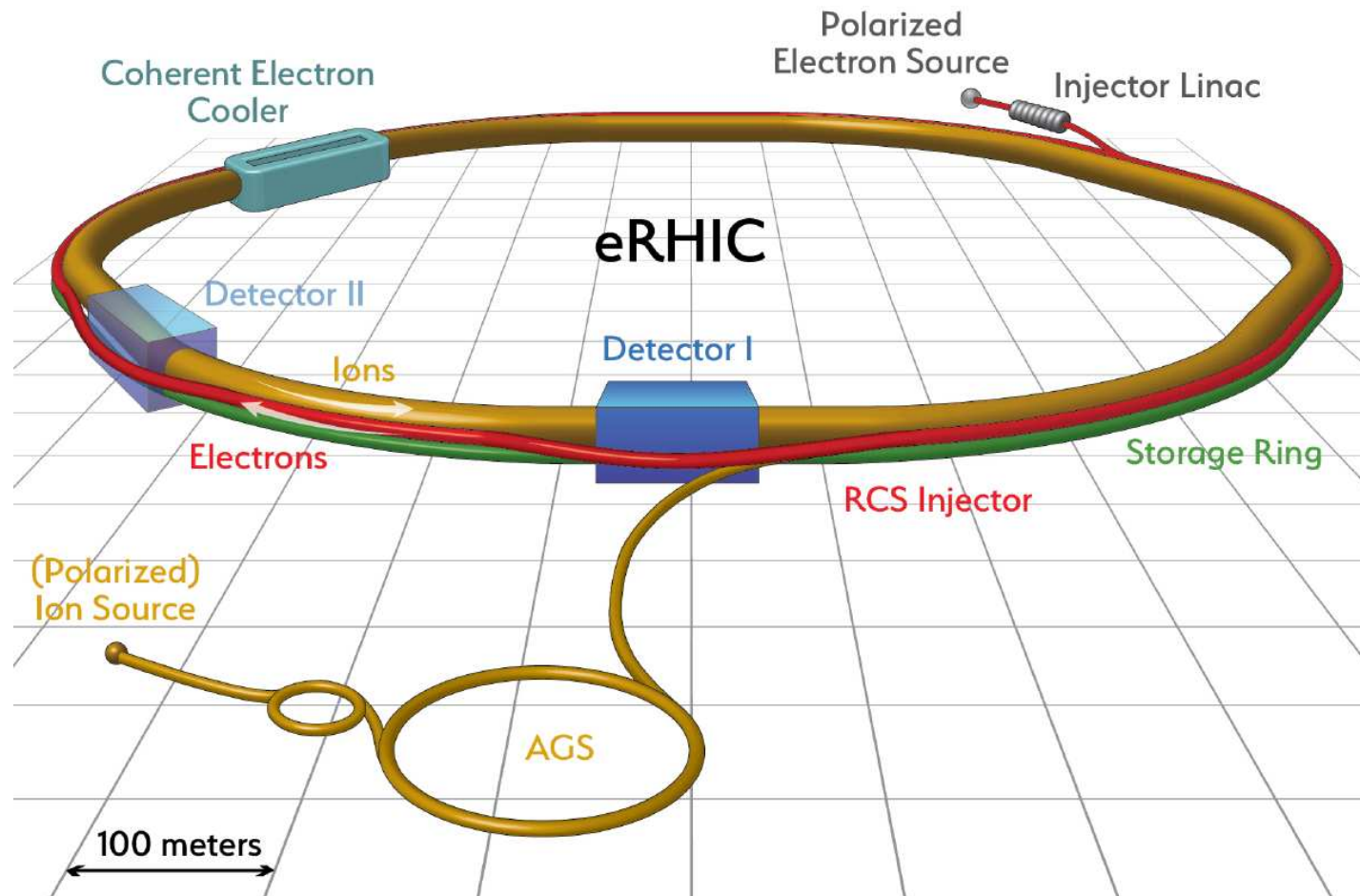


POLARIZATION IN eRHIC ELECTRON STORAGE RING  
AN ERGODIC APPROACH

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# Electron polarization

- From source: 85%, longitudinal
- Into RCS: 85%, vertical, alternately  $\uparrow$  and  $\downarrow$   
*Tilted to vertical at RCS injection*
- RCS transmission to store energy: >95%  
*Provided quadrupole/BPM alignment < 0.2 mm rms, and orbit control < 0.5 mm*
- Bunch into storage ring: >80%,  $\uparrow$  and  $\downarrow$   
*Operation energies: 5, 10 or 18 GeV*
- Stored polarization is affected by synchrotron radiation effects, two processes

# Self-polarization (Sokolov-Ternov effect)

- Synchrotron radiation causes electron spins to slowly self-polarize towards anti- $\vec{B} //$ , equilibrium  $P_{ST}$ ,

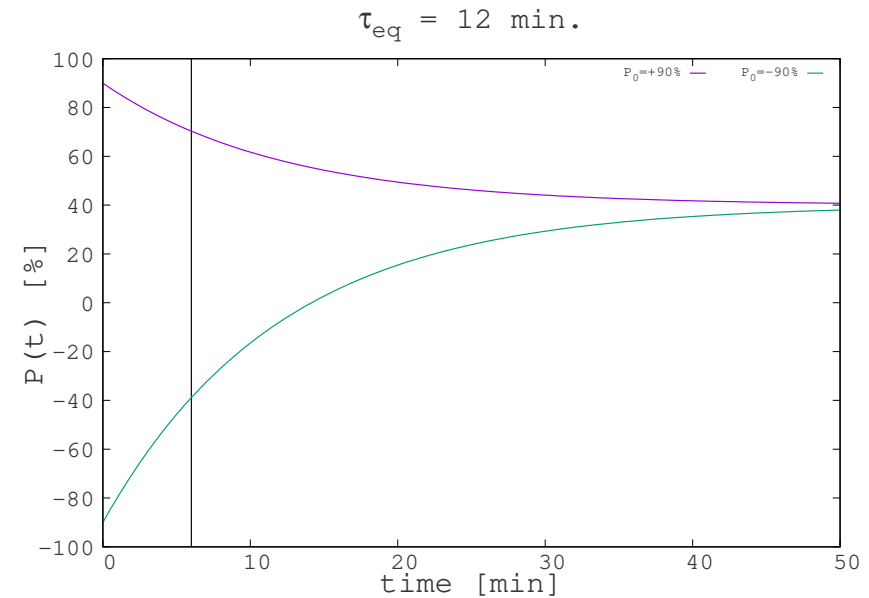
$$P(t) = P_{ST} + (P(0) - P_{ST})e^{-t/\tau_{ST}}$$

$$\diamond \tau_{ST} = \left[ \frac{5\sqrt{3}}{8} \frac{\hbar r_e \gamma^5}{m c} \times \oint_{\text{dipoles}} \frac{ds}{|\rho|^3} \right]^{-1} \xrightarrow{\text{iso-B}} 99 \frac{\rho^2 R}{E_{[GeV]}^5} [\text{sec}]$$

- In particular: bunches injected  $\vec{B} //$  slowly loose their polarization by spin-flip to anti- $\vec{B} //$

- Goal polarization at store :  $\langle 70\% \rangle$

- At 18 GeV, a short  $\tau_{ST} \approx 30$  minute requires replacement of “wrong” polarization bunches every 6 minutes (180 bunches /360)



# Spin diffusion

- Stems from the stochastic change of the spin precession direction due to synchrotron radiation,  $\partial \vec{n}_s / \partial \gamma$

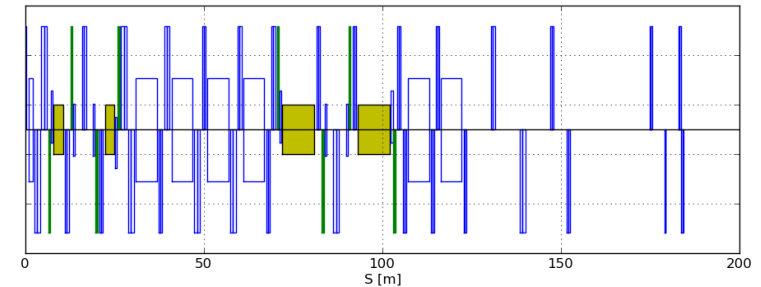
(in a similar way that the chromatic closed-orbit jumps upon SR)

- Dominates bunch depolarization:
  - shortens polarization time,  $\tau_{eq} < \tau_{ST}$ ,
  - reduces asymptotic polarization,  $P_{eq} < P_{ST}$ .

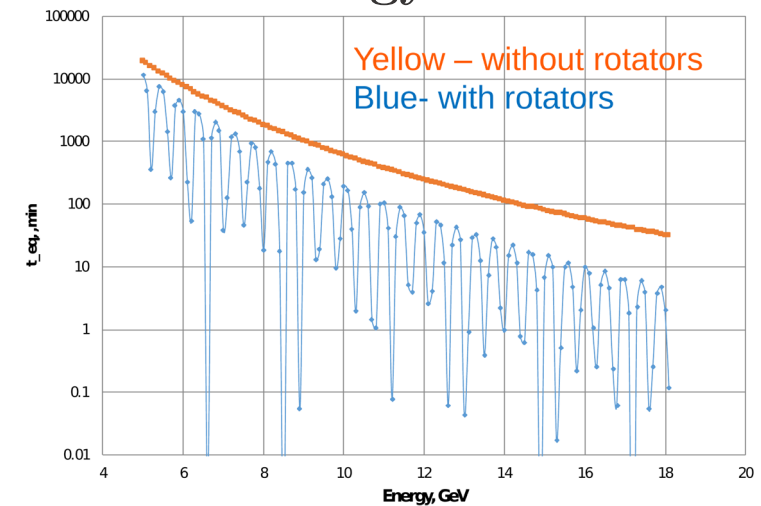
◇ Strength of depolarizing effect  $\propto E^7$

- Spin rotators at eRHIC enhance diffusion as they introduce  
*vertical orbit, vertical dispersion, coupling...*

◇ Two solenoid + two bending magnet sections, interleaved



◇ Equilibrium time  $\tau_{eq}$  versus Energy at eRHIC



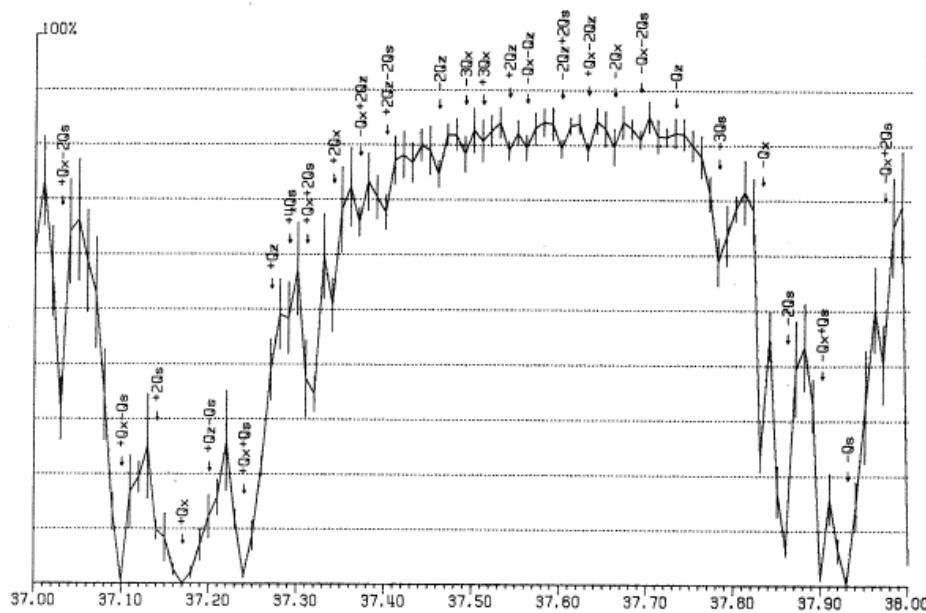
# Assessment of polarization performance

- Typically:

As produced using SITROS at DESY  
*[J. Kewisch, DESY Rep. 83-032, 1983]*

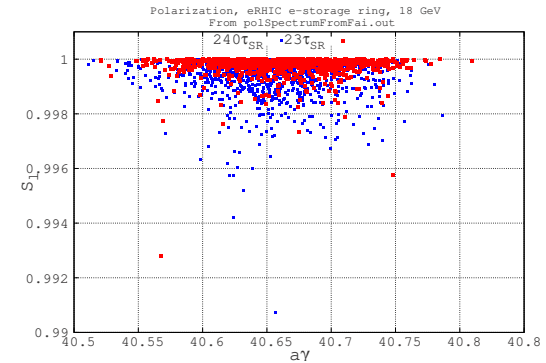
(i) track particles and spins, including Monte Carlo SR,

(ii) produce polarization landscape, *i.e.*,  $P_{eq}$  versus ring rigidity setting ( $a\gamma$  units, here):



HERA-e, linear, res. non-corrected

- Spins tracked over 240 damping times in eRHIC, 18 GeV



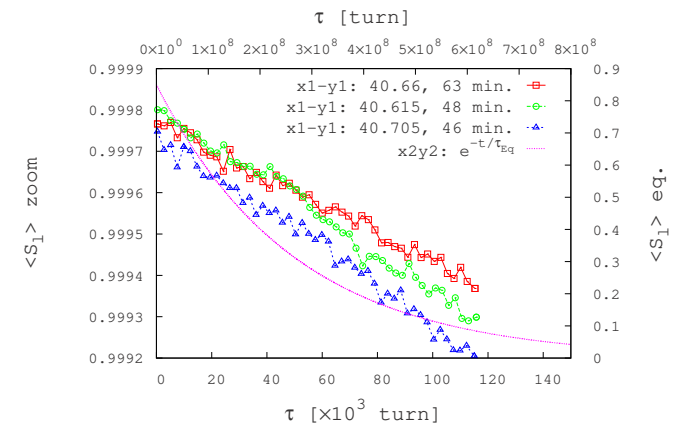
- Polarization decay with time:  
 The diffusion time constant  $\tau_D$  is obtained from linear regression  
 $P/P_0 = \exp(-t/\tau_D) \approx 1 - t/\tau_D$ .

Then,  $P_{eq}$  stems from

$$P_{eq} = P_{ST} \times \tau_{eq} / \tau_{ST}$$

given that  $\tau_{eq} = (1/\tau_{ST} + 1/\tau_D)^{-1}$ .

Polarization vs. time at store, 18 GeV



# This is very much HPC consuming

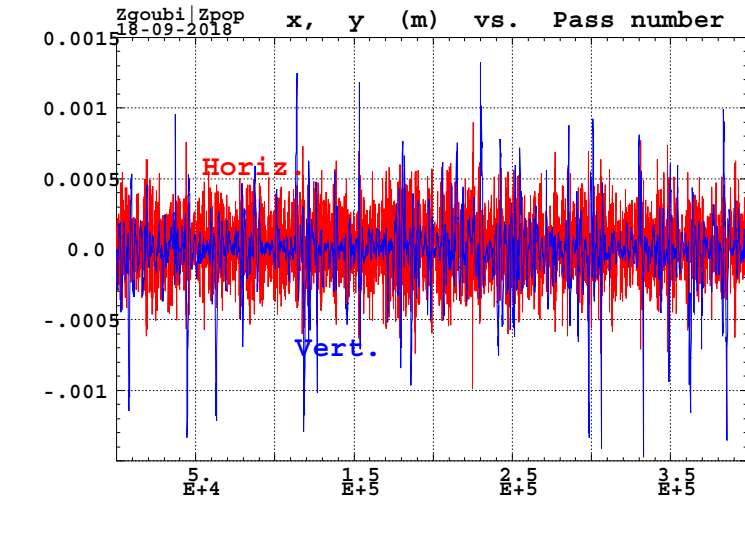
- Typically:
  - ◇ tens of bunches, each its own ring energy
  - ◇ of the order of a thousand particles per bunch
  - ◇ tracking is over several damping times,  
damping time:
    - 500 turns at 18 GeV
    - 3000 turns at 10 GeV
  - ◇ around a large ring - many optical elements (eRHIC circumference 3.833 km).

# Try some ergodic hypotheses, instead

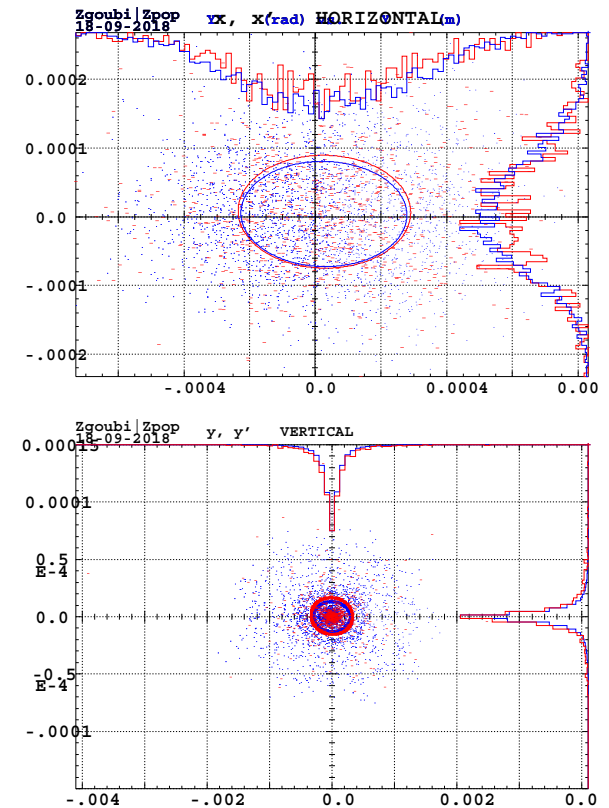
- The dynamical system of an electron bunch in the presence of synchrotron radiation, at equilibrium, is ergodic.

$$\lim_{T \rightarrow \infty} \int_{t_0}^{t_0+T} f(\vec{X}(t)) dt = \int f(\vec{X}) \rho(\vec{X}) d^N \vec{X} \Big|_{\text{time}=t}$$

Transverse particle excursion:



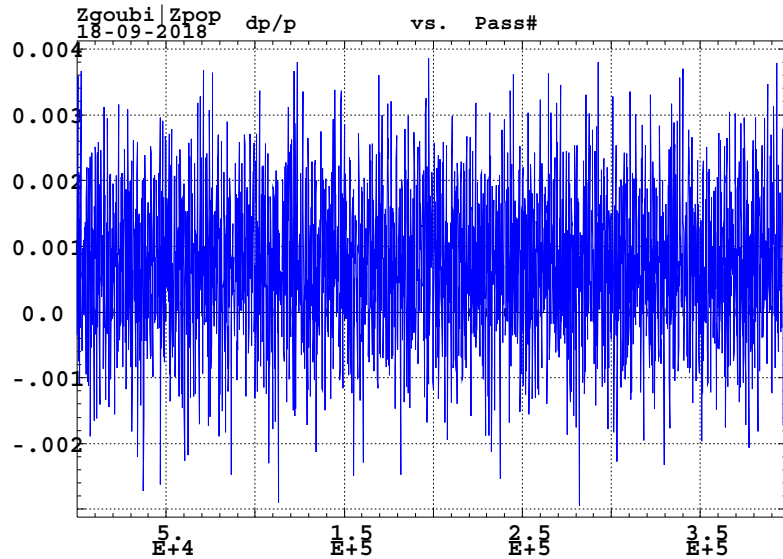
## TRANSVERSE PHASE SPACES.



Blue: single particle motion of the left figure, and matching ellipses. (a 27% coupling, by the solenoid based spin rotator in eRHIC IR6).

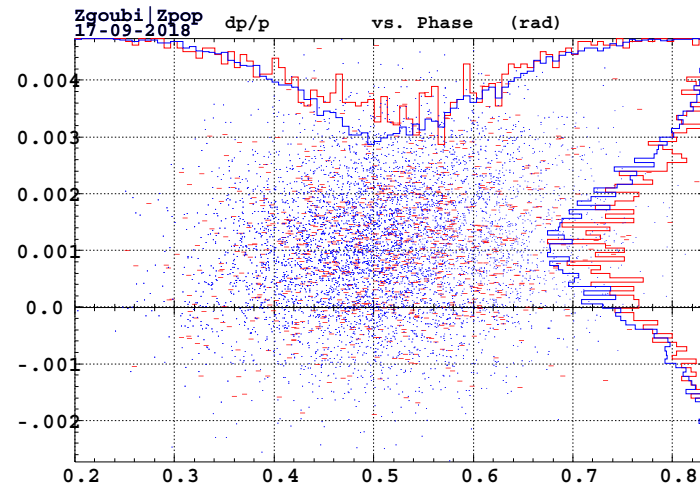
Red: for comparison, case of a  $10^3$  particle bunch, observed at time  $t = 10^3 \tau_{\text{SR}}$ ; rms matching ellipses.

# LONGITUDINAL PHASE SPACE



Stochastic energy excursion over time interval  $t/\tau_{\text{SR}} : 1 \rightarrow 10^3$ .

$\langle \delta p/p \rangle = 1.11 \cdot 10^{-3}$  and  $\sigma_{\delta p/p} = 1.14 \cdot 10^{-3}$ .



**Blue:** projection of the single particle motion of the left plot;

$$\langle \delta p/p \rangle = 1.11 \cdot 10^{-3}$$

$$\sigma_{\delta p/p} = 1.14 \cdot 10^{-3}$$

$$\langle \phi \rangle = 0.519$$

$$\sigma_{\phi} = 0.091;$$

**Red:** a  $10^3$  particle bunch observed at  $t = 10^3 \tau_{\text{SR}}$ ;

$$\langle \delta p/p \rangle = 1.07 \cdot 10^{-3}$$

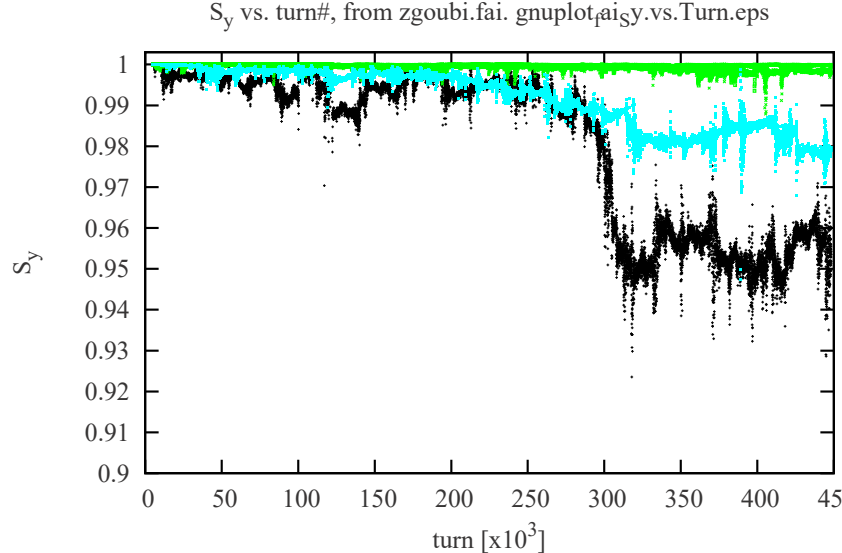
$$\sigma_{\delta p/p} = 1.13 \cdot 10^{-3}$$

$$\langle \phi \rangle = 0.519$$

$$\sigma_{\phi} = 0.091.$$



# Spin motion



- Stochastic spin motion, single particle, observed at IP8.

- Motion is not at equilibrium. However:

- ◊ only ring settings that feature very slow polarization decay are of interest

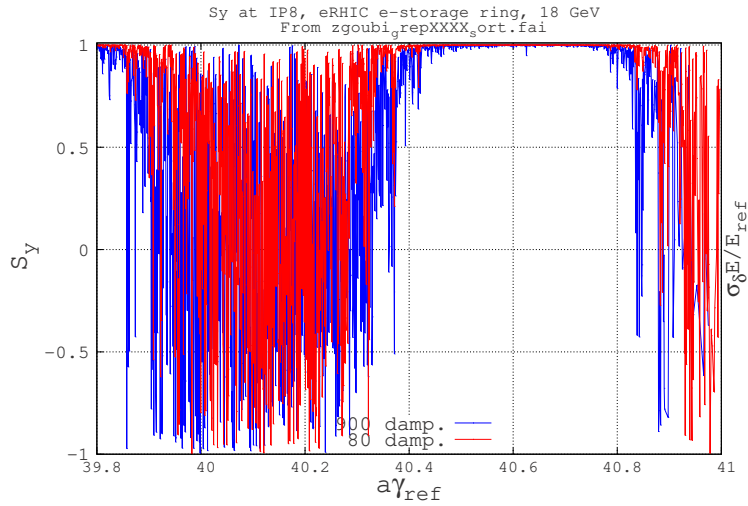
- if the electron motion neighbors depolarizing resonances, these will be revealed by fast decay as  $\tau_D \sim (\alpha\gamma_{Res.} - \alpha\gamma)^2 \tau_{ST}$ ,  $P_{eq} \sim (\alpha\gamma_{Res.} - \alpha\gamma)P_{ST}$ , meaning a ring configuration which is not viable

- In a short time interval, an electron will have explored the all 6D phase-space

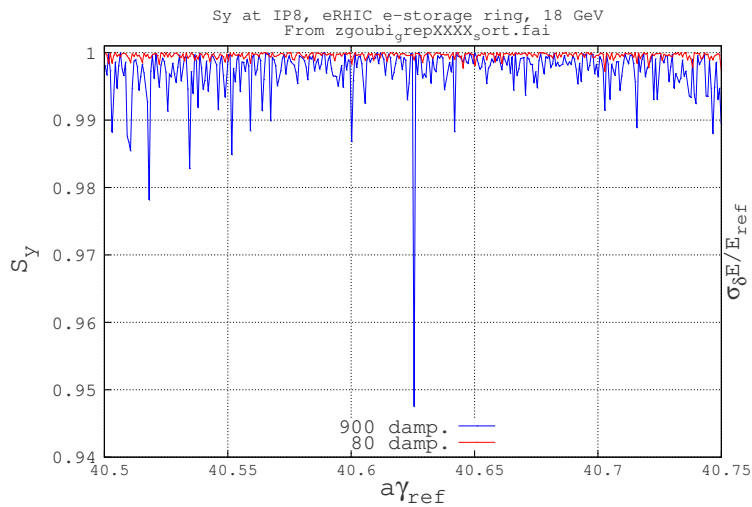
- In a similar way that  $\tau_{SR}$  can be obtained from the observation of the damped motion of a single electron far from equilibrium,  $\tau_D$  can be obtained from long enough observation of spin motion out of equilibrium.

- Track a single particle per bin
- ◊ 1 bin is 1 ring, set for a specific rigidity (here, “ $a\gamma_{\text{ref}}$ ”)
- ◊ about 1000 rings (bins) here

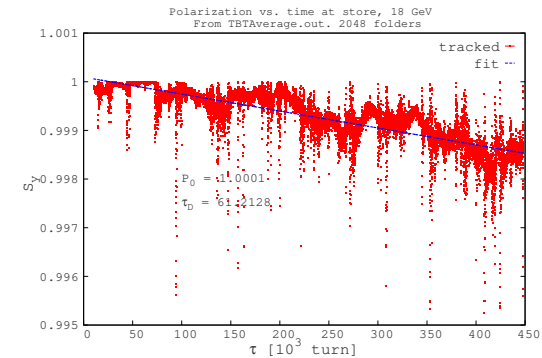
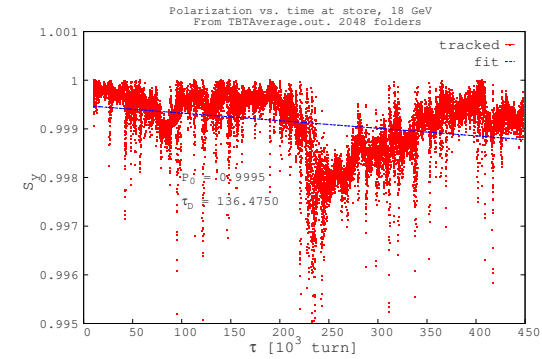
### Spins at 80 and 900 $\times \tau_{\text{SR}}$



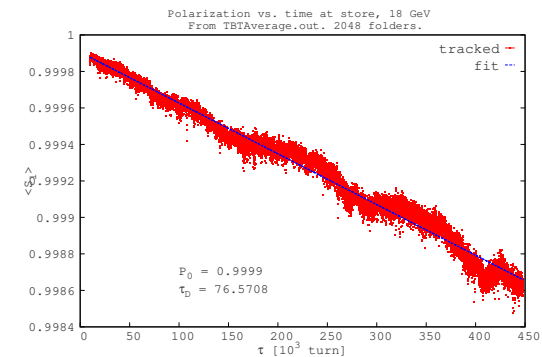
### Zoom-in:



- Monitor individual spins:  
A linear regression on  $P/P_0 = \exp(-t/\tau_D) \approx 1 - t/\tau_D$  provides  $\tau_D$ .

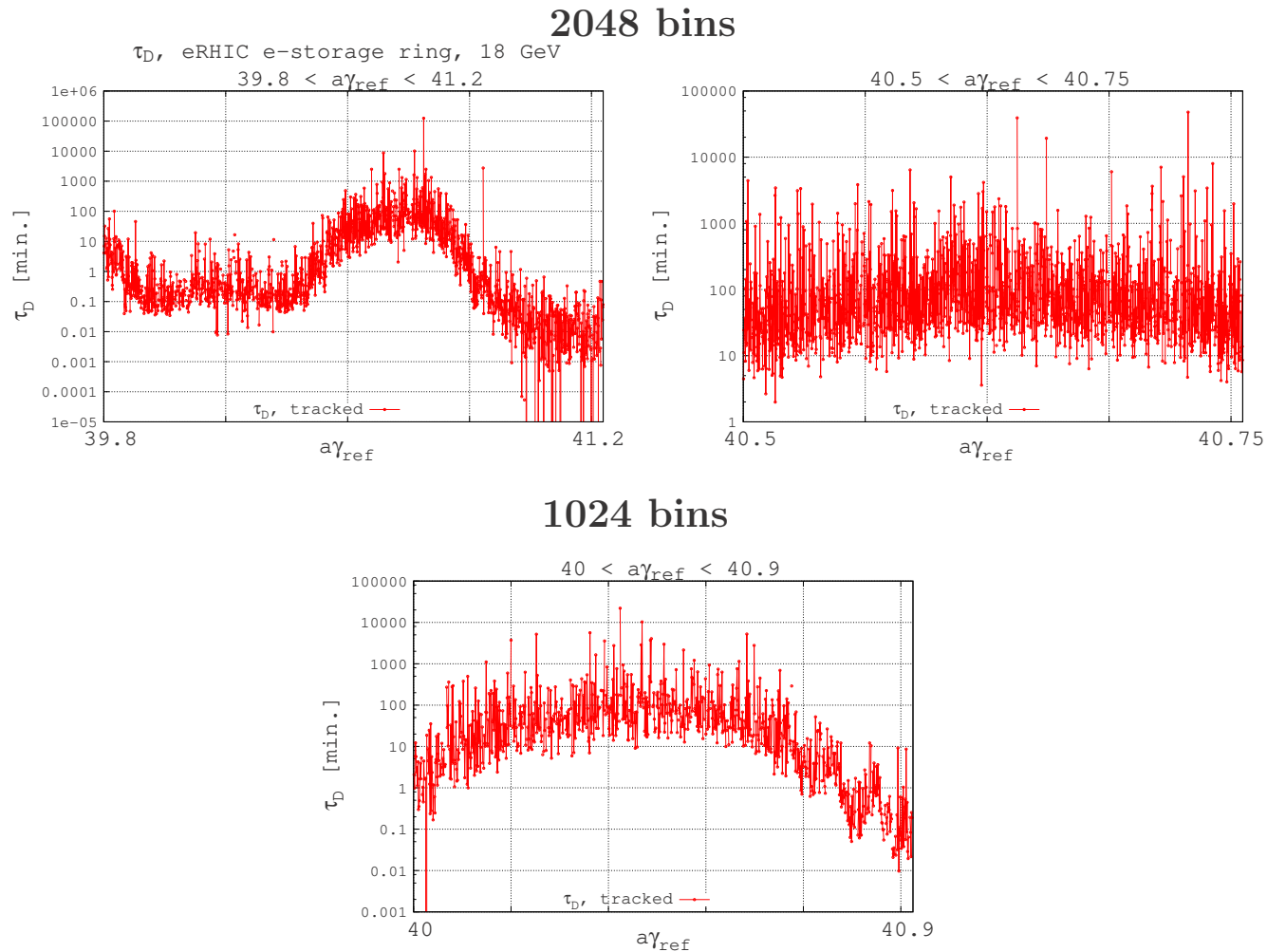


- Possibly, average over reduced  $a\Delta\gamma_{\text{ref}}$  interval, *i.e.*, a few rings/bins ( $\Delta\gamma : 40.60 - 40.62$ , here):

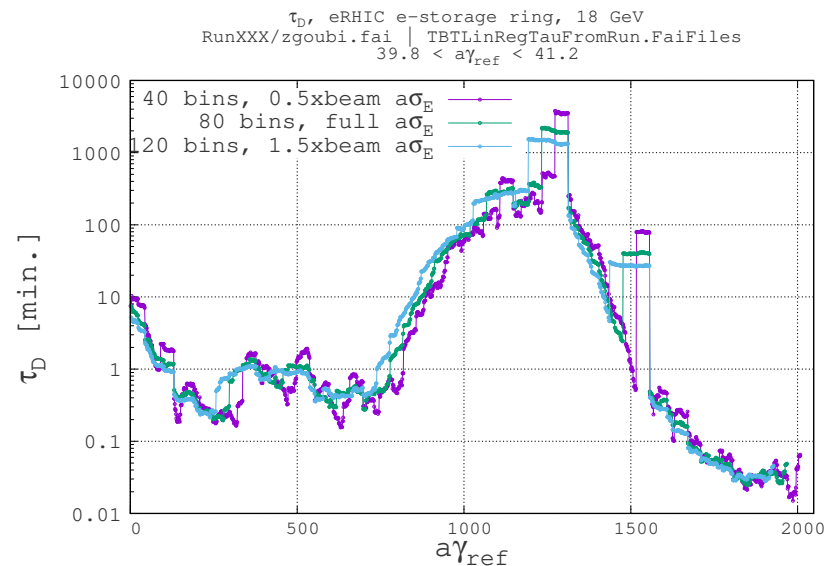
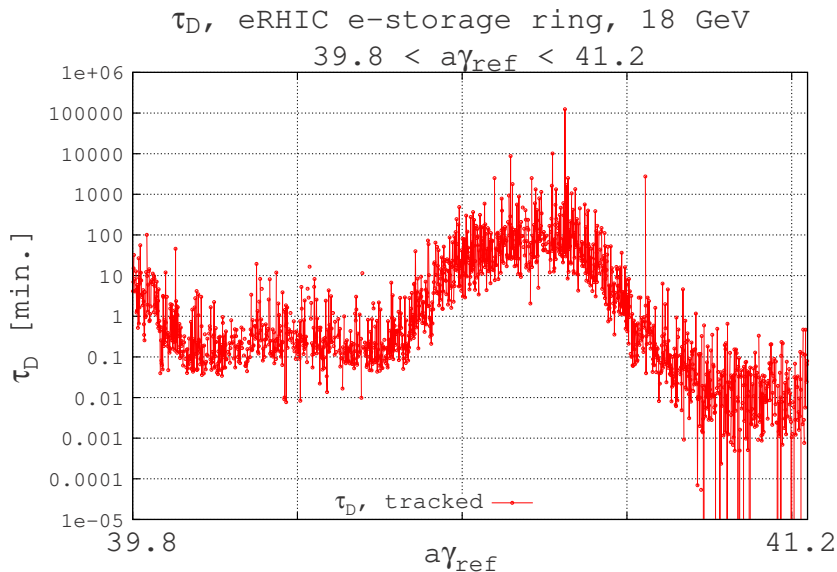


# A metric

- This is in order to compare polarization performance, when optimizing the spin rotator, injecting errors and their compensation, etc.
- A possibility: distance between  $\tau_D$  distributions (or  $P(t)$  distributions at given  $t$ )

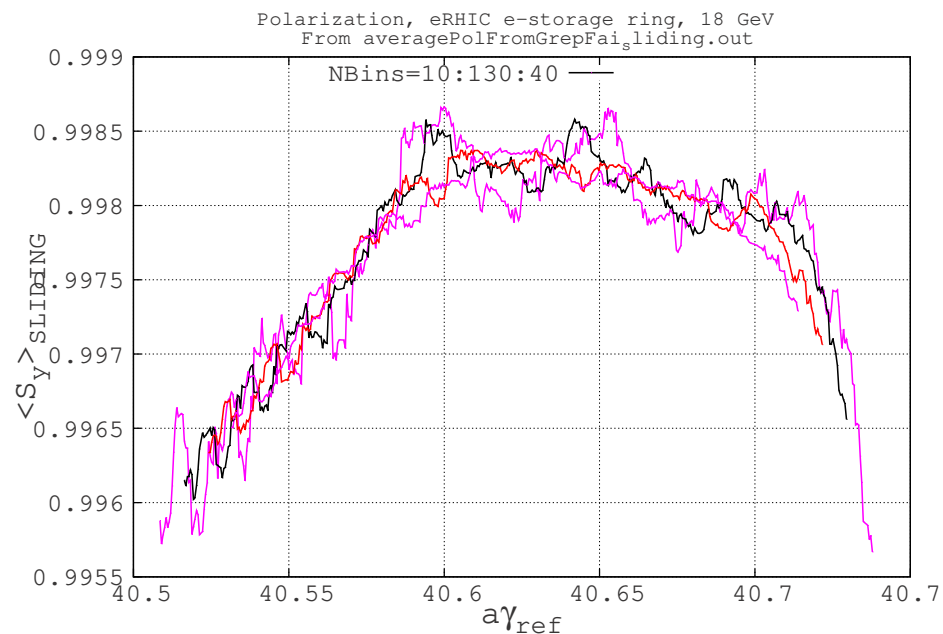
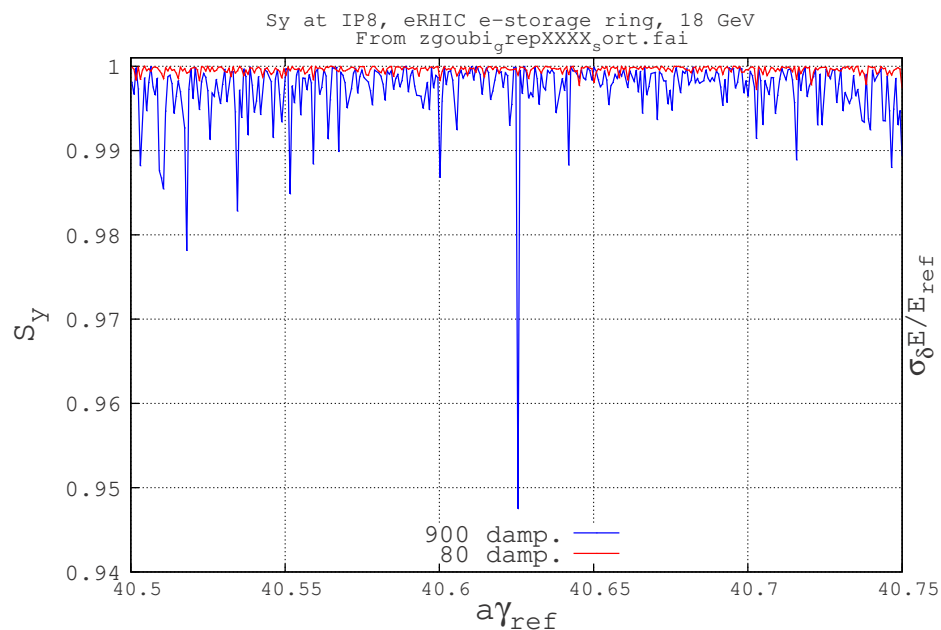


- Averaging over a small  $a\delta\gamma_{\text{ref}}$  interval (a few particles / bins) greatly smooths the fluctuations.
- Justification for a sliding average:
  - (i) with  $a\Delta\gamma_{\text{ref}} = 1$  being covered in  $1024 a\gamma_{\text{ref}}$  bins (or 2048), a beam, which has  $\sigma_\gamma/\gamma_{\text{ref}} \approx 10^{-3}$  (or  $a\delta\gamma_{\text{ref}} = 0.04$ ), extends over about 40 (or 80) bins,
  - (ii) so, a set of a few bins almost belong in the same ring, thus averaging over a few bins is not very different from averaging over a few particles in the same bin
- In this sliding sampling, the distribution appears to evolve only weakly with increasing number of samples, N.



# A metric (cont'd)

- Another possibility for comparing optics:
  - ◇ A sliding average (right plot below) is applied on single particle spin values at a given time out of the multiturn tracking (left plot)
  - ◇ The four curves below differ by the number of bins of the sliding sampling:  $N=10, 50, 90$  and  $130$  bins, respectively.
  - ◇ Again, in this sliding sampling, the distribution appears to evolve only weakly with increasing number of samples,  $N$ .



# COMMENTS

- Assume similar resolution using both methods,
  - “HPC-Hungry” and
  - “Ergodic”,namely, the same number of reference rings, nRings (= number of bins), in a given interval  $a\Delta\gamma_{\text{ref}}$ .  
In the present hypotheses (eRHIC lattice, energy, etc.):
  - first method: the HPC volume is  $n\text{Rings} \times 10^3$  [particles/ring]  $\times$  a few 10s of SR damping times
  - second method: the HPC volume is  $n\text{Rings} \times 1$  to 10 times more SR damping times.This is a two to three orders of magnitude difference.
- Larger HPC volume translates in one or the other of,
  - longer queues, longer computing time, more processors, greater volume of I/Os, larger data analysis HPC volume,...
- Faster computation allows easier exploration of parameter space in design optimizations.
- It remains to determine how close the single-particle method can get to the accuracy of the 1000-particle bunch method (an on-going work).

## CONCLUSION

- The single-particle method seems an efficient first approach for qualifying an evolution of a lattice (optics variants, effects of errors, correction schemes, etc.).
- Plans: use it and improve it at eRHIC!

THANK YOU FOR YOUR ATTENTION

## BIBLIOGRAPHY

- BNL eRHIC collaboration and documents
- eRHIC p-CDR, BNL 205809-2018-FORE (2018)